



Magnetic analysis of 3 and 4-layer dipole magnets

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I. Introduction

Fermilab superconducting magnet group has been working on Nb₃Sn accelerator magnet technology development exploring different superconducting strands and design concepts. Several short dipole models built with the MJR conductor reached the quench currents on a level of 50-60% of the short sample limit. The reason of poor quench performance was identified and associated with the magnetic instabilities in superconducting strands. In order to experimentally confirm the theoretical findings of the magnetic instability model and demonstrate viability of the magnet mechanical structure, a short dipole model based on more stable PIT conductor but identical to the previous models coil geometry and mechanical structure has been fabricated and tested and reached its short sample limit of 10 T after a moderate training.

Moving towards higher fields is the next logical step in the magnet development for exploring the limits of A15 superconductors and addressing the DOE quest for high fields in accelerator-relevant magnets. The PIT conductor would be a preferable choice in terms of the magnet stability; however it lacks the critical current density ($J_c(12T, 4.2K) \sim 2000 \text{ A/mm}^2$) necessary to achieve the high fields and is more sensitive to mechanical stresses than other conductors. The RRP conductor on the other hand offers the highest critical current density ($J_c(12T, 4.2K) \sim 3000 \text{ A/mm}^2$) and lower sensitivity to mechanical stress, but is unstable at low fields. A newly emerging Nb₃Al conductor has critical current density and stability characteristics similar to PIT, but is mechanically superior to any other A15 conductor.

These considerations need to be carefully taken into account during the magnet design process. The purpose of this note is to overview the possible dipole magnet designs that can potentially achieve higher fields based on the present understanding of the conductor behavior and require minimum resources for fabrication at Fermilab by reutilizing as much of the available tooling and equipment as possible.

II. Magnets using HFDA/TQ2a/CC tooling

MJR and RRP are the only Nb₃Sn conductors on the market today able to carry enough current density to achieve high fields at relatively high stresses. Also in part due to this ability they are not stable enough at low fields that was confirmed both theoretically and experimentally on HFDA and HFDM models made of 1-mm MJR and 0.7-mm RRP strands [1]-[3]. Thus to realize the MJR and RRP conductor potential and avoid the stability issues one needs to limit the use of these conductors to the cases where the actual $J_c(B)$ characteristic has a negative slope¹ that corresponds to the field region above $\sim 3 \text{ T}$ [4].

¹ This is a simplified stability criterion that however for these particular conductors is close to the general stability criterion.

It is clear that this requirement can not be fulfilled in a 2-layer shell type magnet with the same cable in the inner and outer layers as there always is a low field region in the outer layer. Thus, for achieving higher fields it is necessary to increase the number of layers. In this regard, the most cost effective and fast way towards higher fields is to choose such layer dimensions that the inner layers fit the HFDA production tooling with the maximum outer radius of 52.0 mm and the outermost layer(s) fit(s) the TQ2a production tooling with the maximum outer radius of 68.2 mm.

a) 3-layer magnet based on 1-mm PIT strand

The first idea was to use the existing HFDA 28-strand cable and add one layer on top of the HFDA coil. To achieve a good field quality and have relatively simple coil geometry it was necessary to re-optimize the inner two layers. However, during the optimization it was possible to remove one wedge from the middle layer while retaining the good field quality. Thus the number of end-parts is the same as in HFDA magnets.

The coil and yoke geometry with the flux density plots are shown in Figures 1-2 and the cable and magnet parameters are summarized in Tables 1-3. The peak field point that defines the magnet performance belongs to the pole turn of the inner layer. If the coil is wound without splices (i.e. the same cable in all the layers) then the only stable enough candidate is the PIT conductor that limits the quench bore field at 13.2 T and the peak coil field at 13.8 T. At this field level the coil stress reach 101 MPa that is comparable with the stresses in HFDA05 magnet.

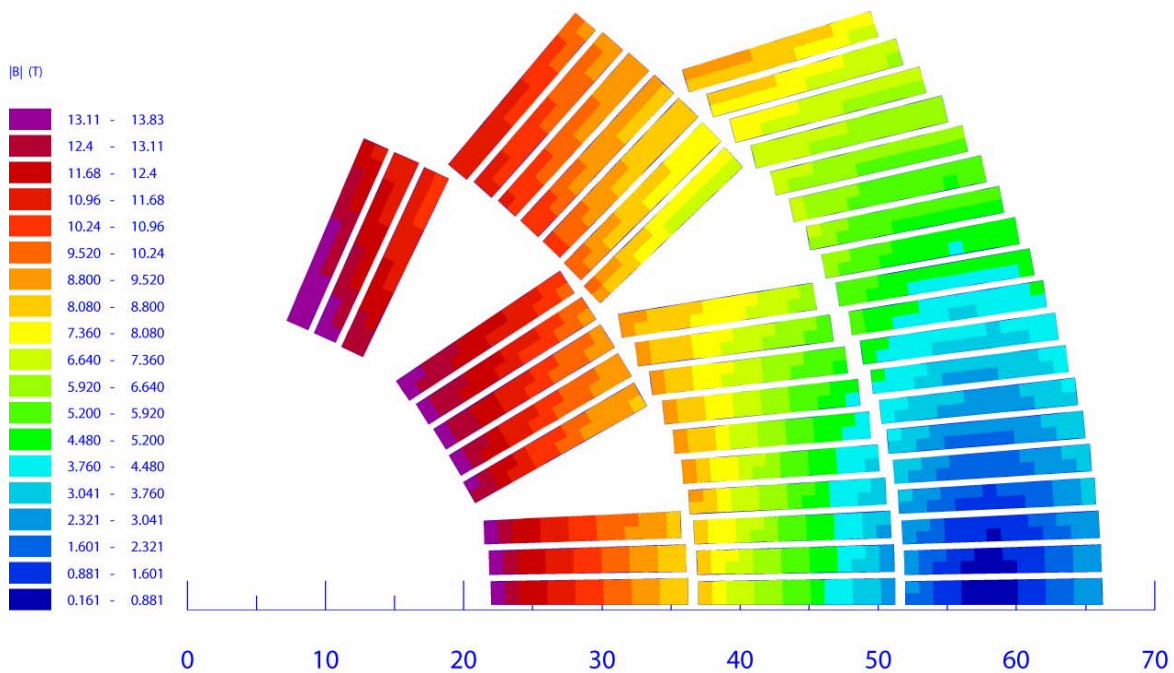


Figure 1. Flux density in the coil cross-section.

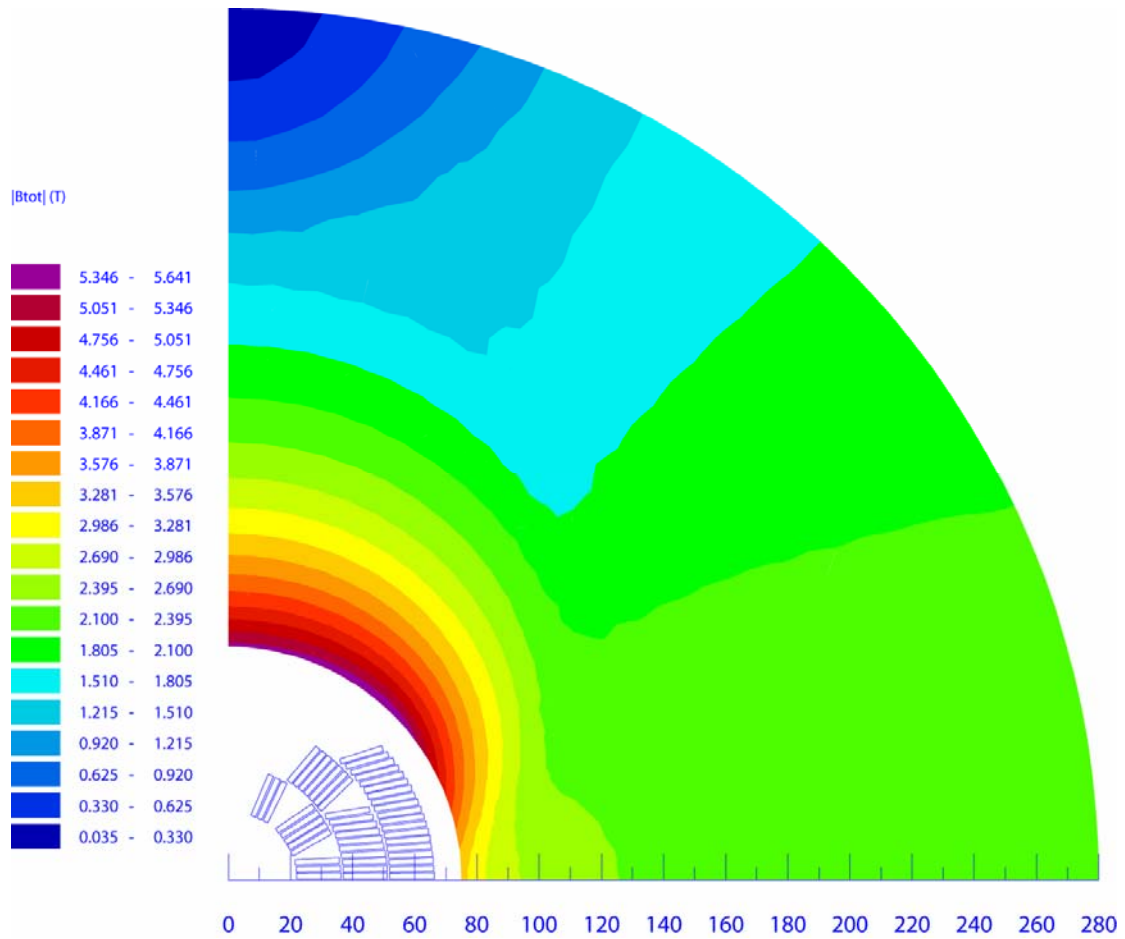


Figure 2. Flux density in the yoke cross-section.

Table 1. Cable parameters.

Parameter	Unit	Inner coil
N of strands	-	28
Strand diameter	mm	1.000
Bare cable width	mm	14.232
Bare inner cable edge thickness	mm	1.687
Bare outer cable edge thickness	mm	1.913
Cabling angle	deg.	14.5
Radial insulation thickness	mm	0.254
Azimuthal insulation thickness	mm	0.254
Copper to non-copper ratio	-	1.000

Table 2. Systematic field harmonics at $R_{\text{bore}}/2$ radius.

n	$b_n, 10^{-4}$	
	I=1kA	I= I_{quench}
3	0.00028	-1.19239
5	0.00098	0.06307
7	0.00299	0.00385
9	-0.01463	-0.01638
11	0.06362	0.07126
13	-0.00236	-0.00265

Table 3. Magnet parameters.

Parameter	Unit	Value
N of layers		3
N of turns in the inner coil		56
N of turns in the outer coil		38
Total coil area (Cu + nonCu)	cm ²	42.70
Assumed non-Cu J_c at 12 T, 4.2 K	A/mm ²	2000
Bore quench field	T	13.24
Quench current	kA	14.86
Peak field in the inner coil at quench	T	13.84
Peak field in the outer coil at quench	T	9.04
Inner coil margin at quench	%	0.00
Outer coil margin at quench	%	26.89
Magnet inductance at quench	mH/m	5.07
Stored energy at quench	kJ/m	559.78
Horizontal Lorentz force/I quadrant at quench	MN/m	4.01
Vertical Lorentz force/I quadrant at quench	MN/m	-1.93
Midplane pressure in the inner coil at quench	MPa	101
Midplane pressure in the outer coil at quench	MPa	79

b) 4-layer magnet based on 1-mm MJR(RRP)/Nb₃Al strands

The maximum field in the magnet design presented in the previous paragraph was limited at 13.2 T field due to the inability to use high- J_c conductors in the coils made of a single piece of cable. Thus a 4-layer design with the different cables in the inner and outer layers was considered.

The coil and yoke geometry with the flux density plots are shown in Figures 3-4 and the cable and magnet parameters are summarized in Tables 4-6. The inner coil is identical to the HFDA type coils. The outer coil is made of twice narrower cable using the same strand that grades the current density between the inner and the outer layers by a factor of two. The quench bore field of 14.1 T is limited by the outer coil layers for the same critical current density in all the layers. As a consequence, this design can not benefit from using RRP conductor in the inner layers (that however is acceptable in terms of the magnet stability), though the quench field is ~1 T higher than in 3-layer design thanks to the grading.

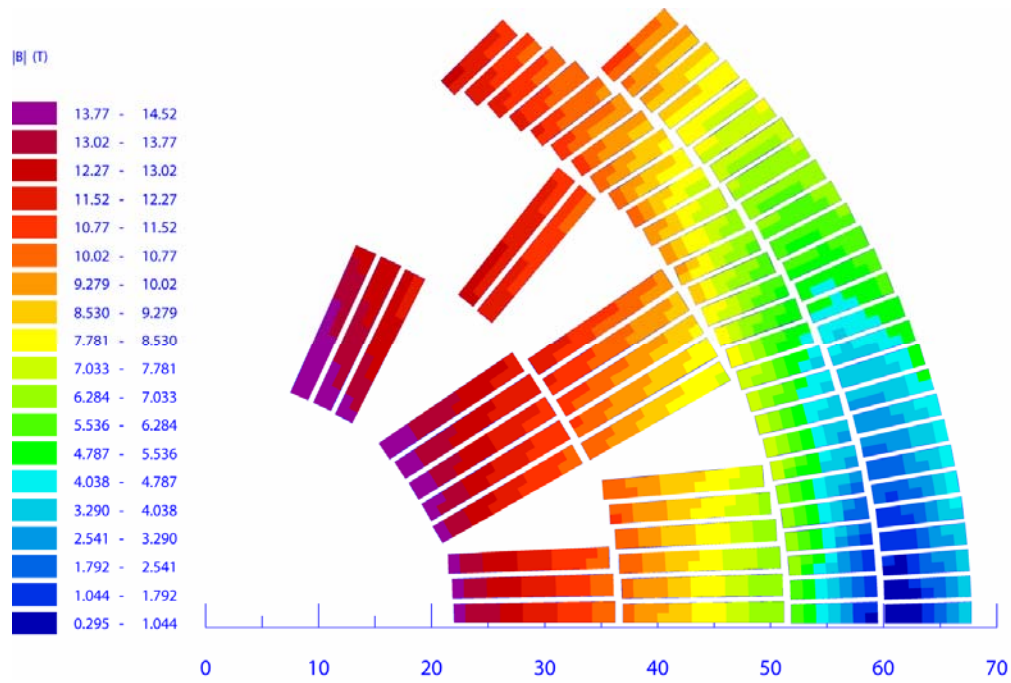


Figure 3. Flux density in the coil cross-section.

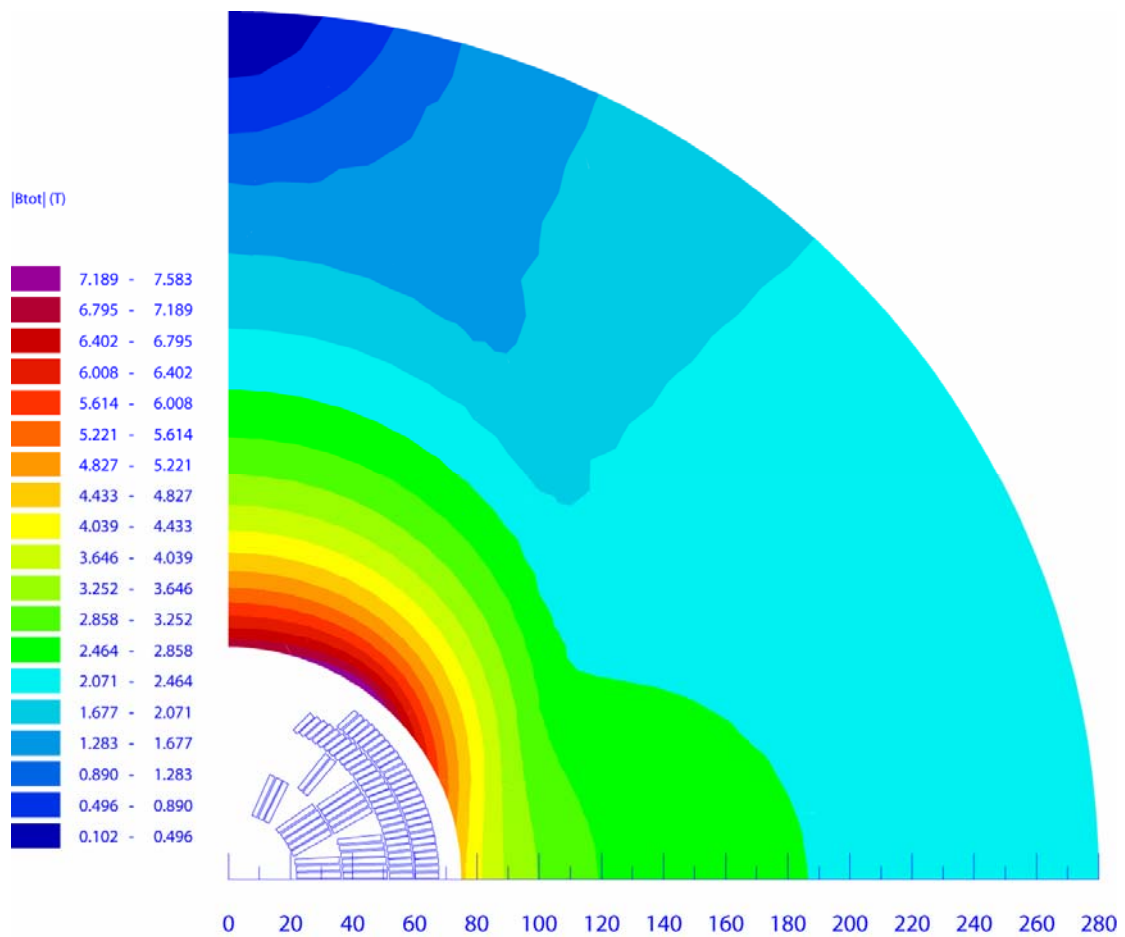


Figure 4. Flux density in the yoke cross-section.

Table 4. Cable parameters.

Parameter	Unit	Inner coil	Outer coil
N of strands	-	28	15
Strand diameter	mm	1.000	1.000
Bare cable width	mm	14.232	7.635
Bare inner cable edge thickness	mm	1.687	1.669
Bare outer cable edge thickness	mm	1.913	1.882
Cabling angle	deg.	14.5	14.5
Radial insulation thickness	mm	0.254	0.225
Azimuthal insulation thickness	mm	0.254	0.225
Copper to non-copper ratio	-	1.000	1.000

Table 5. Systematic field harmonics at $R_{\text{bore}}/2$ radius.

n	$b_n, 10^{-4}$	
	I=1kA	I= I_{quench}
3	-1.0754	0.13626
5	-0.6008	-0.63001
7	0.0122	0.01443
9	-0.0359	-0.04327
11	0.0391	0.04706
13	-0.0007	-0.00081

Table 6. Magnet parameters.

Parameter	Unit	Value
N of layers		4
N of turns in the inner coil		48
N of turns in the outer coil		110
Total coil area (Cu + nonCu)	cm ²	48.58
Assumed $J_c^{\text{non-Cu}}$ (12 T, 4.2 K)	A/mm ²	2000
Bore quench field	T	14.08
Quench current	kA	10.92
Peak field in the inner coil at quench	T	15.06
Peak field in the outer coil at quench	T	12.37
Inner coil margin at quench	%	3.62
Outer coil margin at quench	%	0.00
Magnet inductance at quench	mH/m	13.99
Stored energy at quench	kJ/m	834.13
Horizontal Lorentz force/I quadrant at quench	MN/m	5.21
Vertical Lorentz force/I quadrant at quench	MN/m	-2.47
Midplane pressure in the inner coil at quench	MPa	61.6
Midplane pressure in the outer coil at quench	MPa	242.6

Due to the twice narrower cable, the midplane pressure in the outer layers reach a whopping 243 MPa that is way above the tolerable pressure for Nb₃Sn conductor thus leaving the Nb₃Al conductor as the only feasible option for the outer layers.

c) 4-layer magnet based on 1-mm MJR(RRP)/PIT strands

To accommodate the more pressure-sensitive Nb₃Sn conductors for the 4-layer magnet design presented earlier, a number of turns was removed from the outer layers.

The coil and yoke geometry with the flux density plots are shown in Figures 5-6 and the cable and magnet parameters are summarized in Tables 7-9. Removing the turns heavily impacted the field quality with 65 units of the sextupole component as well as the bore quench field that dropped to 13.5 T – that is just 0.3 T higher than the quench field of the 3-layer magnet.

Similarly to the previous design, the magnet performance is limited by the outer layers for the same critical current density in all the layers. The midplane pressure was reduced below the “magic” for a Nb₃Sn conductor value of 150 MPa, though may still be too large for the highly pressure-sensitive PIT conductor.

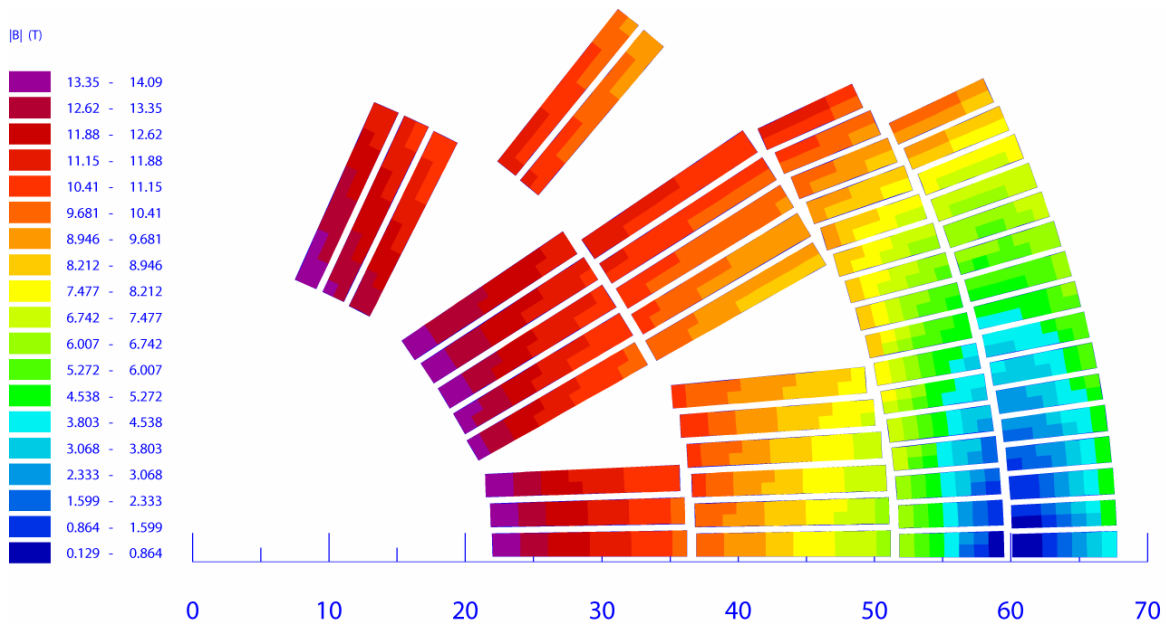


Figure 5. Flux density in the coil cross-section.

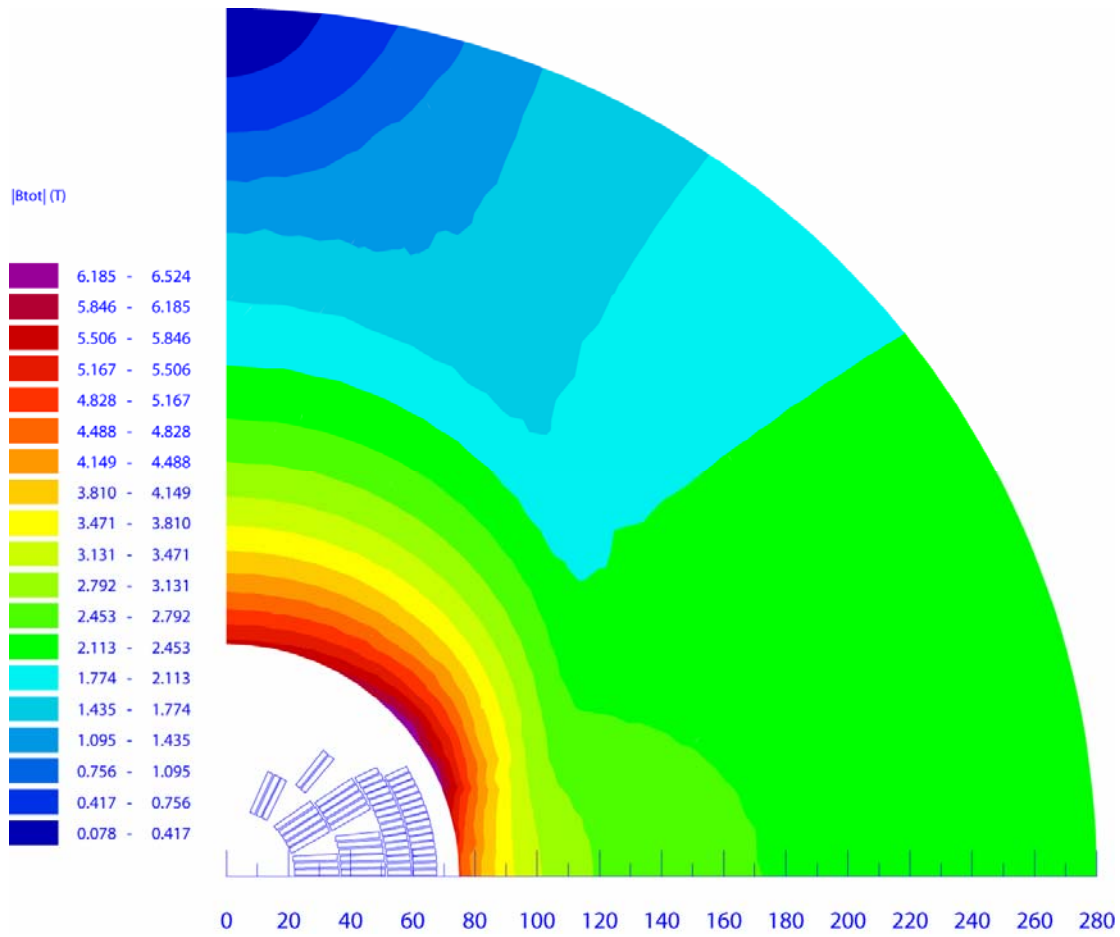


Figure 6. Flux density in the yoke cross-section.

Table 7. Cable parameters.

Parameter	Unit	Inner coil	Outer coil
N of strands	-	28	15
Strand diameter	mm	1.000	1.000
Bare cable width	mm	14.232	7.635
Bare inner cable edge thickness	mm	1.687	1.669
Bare outer cable edge thickness	mm	1.913	1.882
Cabling angle	deg.	14.5	14.5
Radial insulation thickness	mm	0.254	0.225
Azimuthal insulation thickness	mm	0.254	0.225
Copper to non-copper ratio	-	1.000	1.000

Table 8. Systematic field harmonics at $R_{\text{bore}}/2$ radius.

n	$b_n, 10^{-4}$	
	I=1kA	I= I_{quench}
3	65.3208	65.2674
5	0.0868	0.1083
7	-0.0194	-0.0217
9	-0.0460	-0.0538
11	0.0500	0.0584
13	-0.0009	-0.0010

Table 9. Magnet parameters.

Parameter	Unit	Value
N of layers		3
N of turns in the inner coil		48
N of turns in the outer coil		64
Total coil area (Cu + nonCu)	cm ²	37.38
Assumed non-Cu J_c at 12 T, 4.2 K	A/mm ²	2000
Bore quench field	T	13.53
Quench current	kA	13.02
Peak field in the inner coil at quench	T	14.35
Peak field in the outer coil at quench	T	11.47
Inner coil margin at quench	%	1.82
Outer coil margin at quench	%	0.00
Magnet inductance at quench	mH/m	8.33
Stored energy at quench	kJ/m	706.05
Horizontal Lorentz force/I quadrant at quench	MN/m	4.17
Vertical Lorentz force/I quadrant at quench	MN/m	-2.50
Midplane pressure in the inner coil at quench	MPa	74
Midplane pressure in the outer coil at quench	MPa	146

III. Magnets partially using HFDA/TQ2a/CC tooling

The magnet designs presented in the previous section made the maximum use of the available tooling. However, it is clear that these designs can not fully unfold the potential of the RRP conductor as with the tooling-defined layer radii they are limited by the conductor performance in the outer layer(s). Thus a redistribution of the cable width was necessary to charge the inner coil layers with the RRP current density.

a) 4-layer magnet fitting TQ2a tooling based on 1-mm RRP/PIT(Nb_3Al) strands

The coil aperture of this magnet was chosen the same as in HFDA magnets and the coil outer diameter was chosen to fit the TQ2a tooling. The cable size was optimized to fill the space between these two constrains in a 4-layer configuration. In this case at least the HFDA winding mandrel and TQ2a curing, reaction and impregnation fixtures can be used without changes. The coil and yoke geometry with the flux density plots are shown in Figures 7-8 and the cable and magnet parameters are summarized in Tables 10-12.

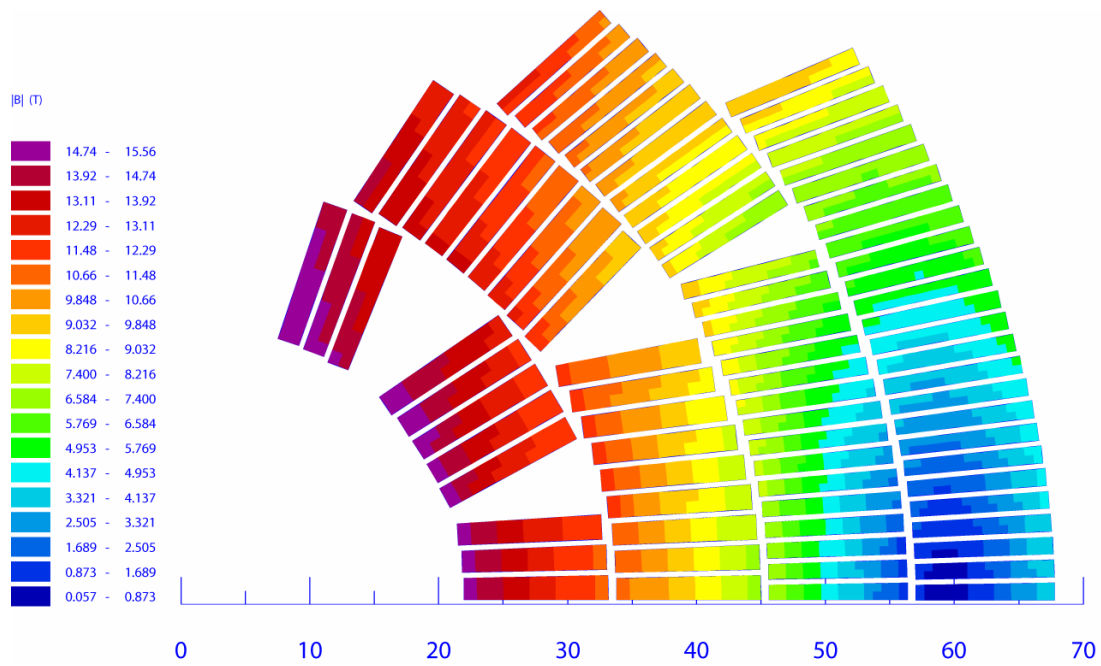


Figure 7. Flux density in the coil cross-section.

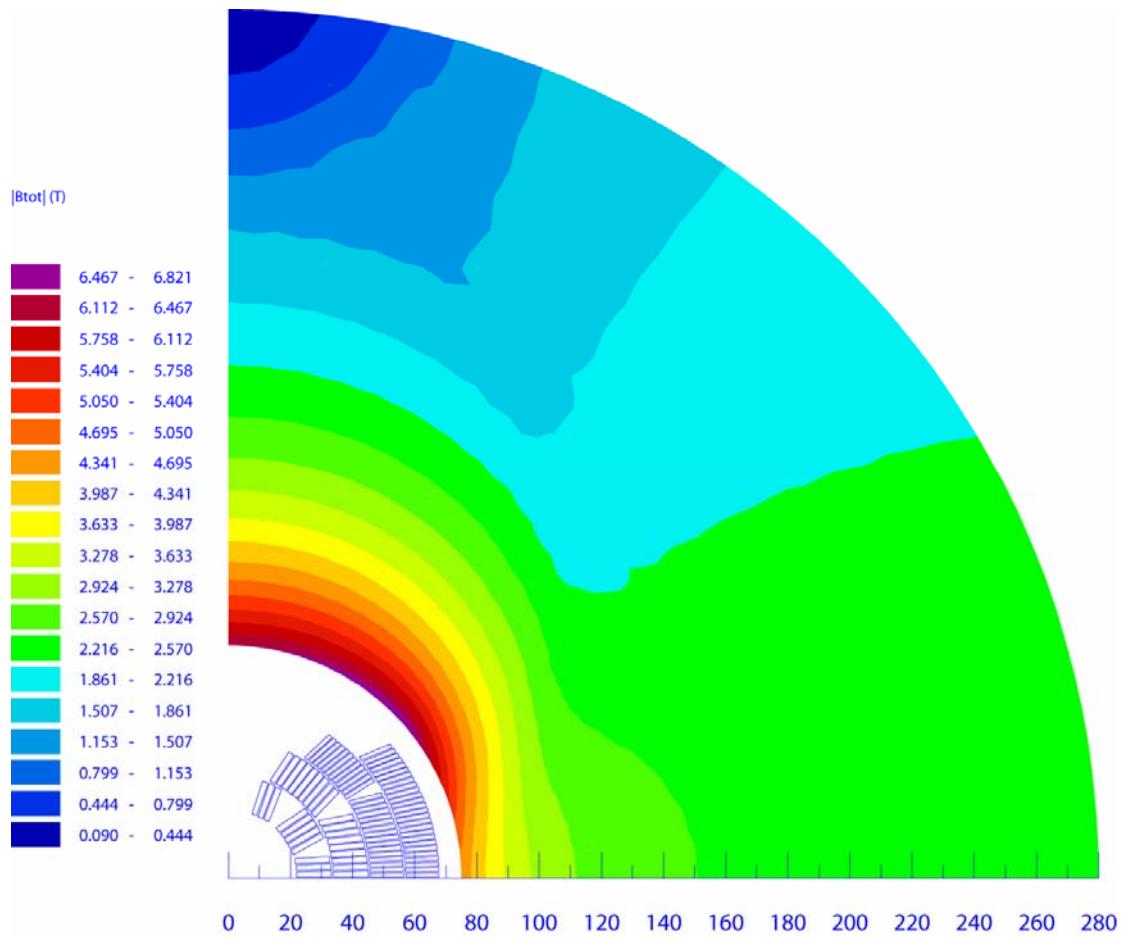


Figure 8. Flux density in the yoke cross-section.

Table 10. Cable parameters.

Parameter	Unit	Inner coil	Outer coil
N of strands	-	22	30
Strand diameter	mm	1.000	0.700
Bare cable width	mm	11.183	10.742
Bare inner cable edge thickness	mm	1.676	1.165
Bare outer cable edge thickness	mm	1.910	1.333
Cabling angle	deg.	14.5	14.5
Radial insulation thickness	mm	0.225	0.225
Azimuthal insulation thickness	mm	0.225	0.225
Copper to non-copper ratio	-	1.000	1.000

Table 11. Systematic field harmonics at $R_{\text{bore}}/2$ radius.

n	$b_n, 10^{-4}$	
	I=1kA	I=I _{quench}
3	0.0003	-3.6583
5	-0.0008	0.0556
7	-0.0043	-0.0046
9	-0.0307	-0.0360
11	0.0489	0.0574
13	-0.0011	-0.0013

Table 12. Magnet parameters.

Parameter	Unit	Value
N of layers		4
N of turns in the inner coil		58
N of turns in the outer coil		108
Total coil area (Cu + nonCu)	cm ²	44.46
Assumed non-Cu J_c at 12 T, 4.2 K in the inner coil	A/mm ²	2800
Assumed non-Cu J_c at 12 T, 4.2 K in the outer coil	A/mm ²	2000
Bore quench field	T	14.98
Quench current	kA	10.27
Peak field in the inner coil at quench	T	15.81
Peak field in the outer coil at quench	T	12.58
Inner coil margin at quench	%	1.59
Outer coil margin at quench	%	0.00
Magnet inductance at quench	mH/m	15.54
Stored energy at quench	kJ/m	819.52
Horizontal Lorentz force/I quadrant at quench	MN/m	5.39
Vertical Lorentz force/I quadrant at quench	MN/m	-2.70
Midplane pressure in the inner coil at quench	MPa	156.6
Midplane pressure in the outer coil at quench	MPa	125.7

Due to the cable width redistribution it is possible to use (and benefit from) the RRP conductor in the inner layers while using the lower-Jc PIT or Nb₃Al conductors in the outer layers. With these conductors the magnet would reach the bore quench field of 15.0 T, while the midplane pressures are within reasonable limits of 157 MPa for the inner layers and 126 MPa for the outer layers.

It is interesting to compare parameters of this magnet with the parameters of 15-T block type magnet based on the RRP conductor with the same characteristics [5]. While the coil of the block type design may look more compacted due to absence of the angle-compensating wedges, it reaches the same quench field in 24% smaller useful aperture at 31 % larger conductor area and 15 % higher stresses under full excitation. Moreover, since the block-type coil uses the same high-Jc strand everywhere, including the low-field regions, this magnet is potentially unstable that depending on the strand RRR may or may not reach the designed field. This comparison clearly demonstrates the advantages of the shell-type coils over the block-type designs.

b) 4-layer magnet not-fitting TQ2a tooling based on 1-mm RRP/PIT(Nb₃Al) strands

The 4-layer magnet design presented in the previous paragraph was optimized to fit the TQ2a tooling that required new (or modified) winding and curing tooling for all but the outermost layer. Another possibility is to use the unchanged HFDA winding, curing, reaction and impregnation tooling for the inner layers similarly to the 4-layer designs described in the previous section, but add the outer layers of the same thickness. Obviously, it will require all new tooling for the outer layers.

The coil and yoke geometry with the flux density plots are shown in Figures 9-10 and the cable and magnet parameters are summarized in Tables 13-15.

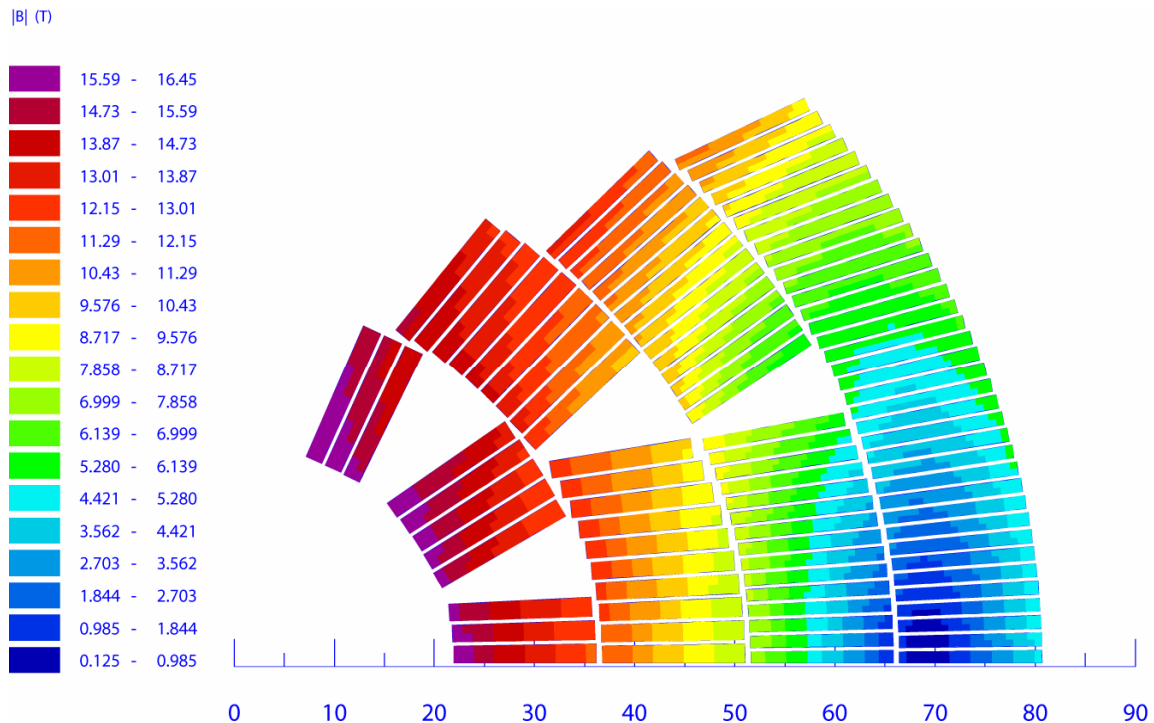


Figure 9. Flux density in the coil cross-section.

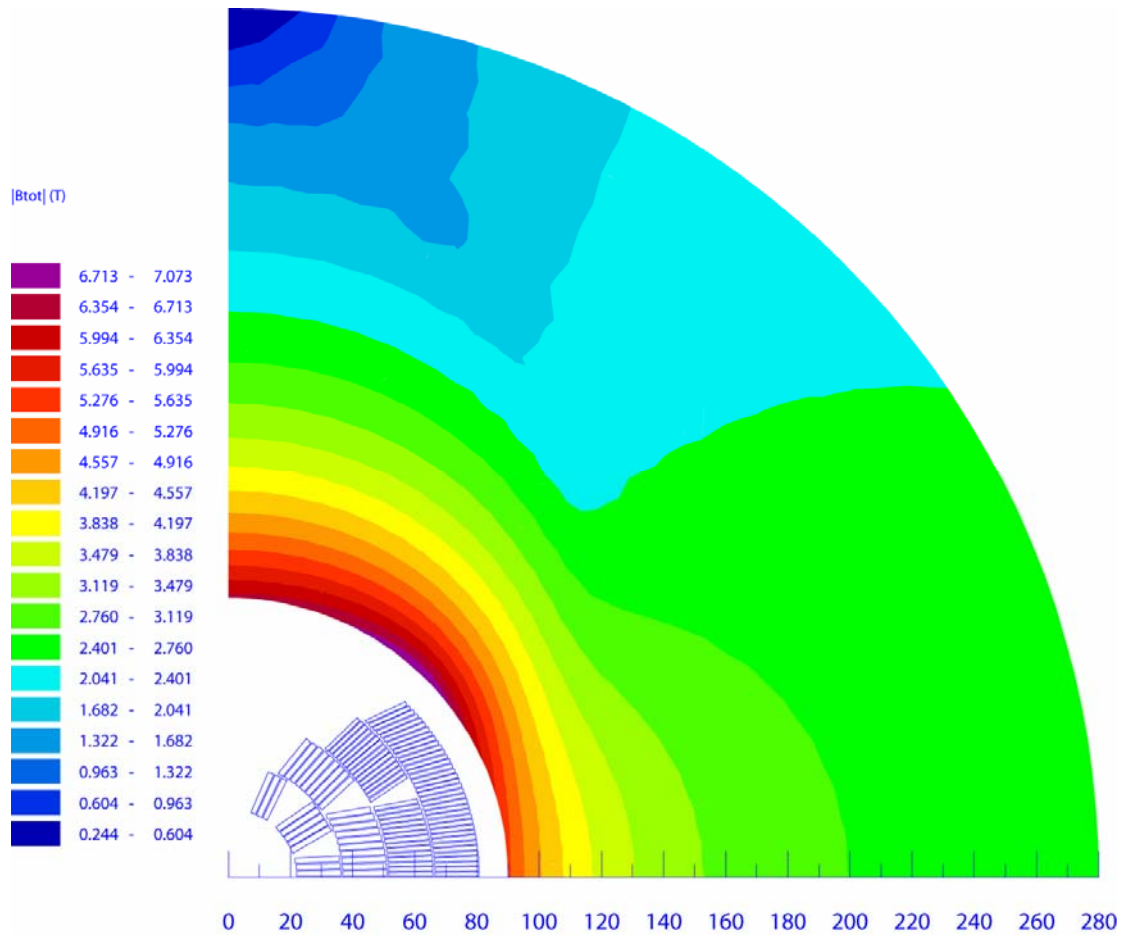


Figure 10. Flux density in the yoke cross-section.

Table 13. Cable parameters.

Parameter	Unit	Inner coil	Outer coil
N of strands	-	28	39
Strand diameter	mm	1.000	0.700
Bare cable width	mm	14.232	14.232
Bare inner cable edge thickness	mm	1.687	1.140
Bare outer cable edge thickness	mm	1.913	1.314
Cabling angle	deg.	14.5	14.5
Radial insulation thickness	mm	0.200	0.200
Azimuthal insulation thickness	mm	0.200	0.200
Copper to non-copper ratio	-	1.000	1.000

Table 14. Systematic field harmonics at $R_{\text{bore}}/2$ radius.

n	$b_n, 10^{-4}$	
	I=1kA	I= I_{quench}
3	0.0001	-3.3333
5	0.0002	0.0273
7	0.0019	0.0024
9	0.0022	0.0026
11	0.0441	0.0526
13	-0.0013	-0.0015

Table 15. Magnet parameters

Parameter	Unit	Value
N of layers		4
N of turns in the inner coil		60
N of turns in the outer coil		132
Total coil area (Cu + nonCu)	cm ²	68.18
Assumed non-Cu J_c at 12 T, 4.2 K in the inner coil	A/mm ²	2800
Assumed non-Cu J_c at 12 T, 4.2 K in the outer coil	A/mm ²	2000
Bore quench field	T	15.96
Quench current	kA	11.22
Peak field in the inner coil at quench	T	16.45
Peak field in the outer coil at quench	T	13.37
Inner coil margin at quench	%	0.00
Outer coil margin at quench	%	0.90
Magnet inductance at quench	mH/m	20.04
Stored energy at quench	kJ/m	1261.40
Horizontal Lorentz force/I quadrant at quench	MN/m	6.74
Vertical Lorentz force/I quadrant at quench	MN/m	-3.57
Midplane pressure in the inner coil at quench	MPa	111.8
Midplane pressure in the outer coil at quench	MPa	143.0

Similarly to the previous design, it was possible to grade the inner and outer layers such that the inner layers use the maximum performance of the RRP conductor and the outer layers use stable PIT or Nb₃Al conductor. The magnet reaches the quench bore field of 16.0 T at a moderate midplane pressure of 112 MPa in the inner layers and high but tolerable pressure of 143 MPa in the outer layers.

IV. CONCLUSION

Different magnetic designs of the shell-type magnets that can be made using available HFDA/TQ2a/CC tooling were developed. Depending on the amount of available tooling used and possibility of modifying old or acquiring new tooling, the magnets can reach 13.2-16.0 T fields.

It was shown that in the case of complete adoption of HFDA/TQ2a tooling the magnets are limited by the conductor performance in the outer layers and can reach a maximum of 13.5 T field.

Dropping the requirement of the intermediate layers to match the existing tooling offers more flexibility in choosing the layer dimensions that is beneficial for achieving fields up to 16.0 T.

A comparison between shell and block-type coils with the same quench field demonstrated that the shell-type coil is more effective than the block-type in terms of the coil aperture, conductor area and maximum stress under excitation.

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